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Optimizing the Characteristics of the Motion of Steel Balls and their Impact on Shell Liners in SAG Mills

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The equations governing the motion of steel balls and their impact onto shell liners in industrial Semi-Autogenous Grinding (SAG) mills are derived in full details by the authors and are used in order to determine the effective design variables for optimizing the working conditions of the mill and to avoid severe impacts which lead to the breakage of SAG mill shell liners. These design variables are the lifter height H , the working coefficient of friction μ , lifter face inclination angle φ , steel ball size r_B , mill rotational velocity ω , and mill size R . In order to optimize the operating conditions and avoid severe impacts to its shell liners, the effect of these parameters need to be studied. The effect of lifter height (H) and the coefficient of friction (μ) on some of the main impact characteristics are simultaneously investigated for a SAG mill in Sarcheshmeh Copper Complex. It was shown that as the lifter height or the coefficient of friction increases, the impact position tends to move upward, but the maximum impact force and the absolute value of the maximum principal stress decreases, reducing the impact severity.

Keywords: Comminution, Semi-Autogenous Grinding Mill (SAG Mill), Liner, Impact, Hertz Contact Law.

1- Introduction

Tumbling mills, which are a critical part of all mining operations and have been used for over 100 years to transform masses of ores to usable-size particles, consist of a cylindrical chamber filled with balls and/or rocks rotating about their longitudinal axes, as shown in Figure 1.

Autogenous grinding (AG) is the action of rocks grinding upon themselves. Grinding in AG mills is frequently assisted by steel balls, in which case it is called a Semi-Autogenous Grinding (SAG) mill. Semi-Autogenous Grinding mills came into operation in the mid 1970s [1].

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Breakage in tumbling mills is grouped into impact breakage, abrasion breakage and attrition breakage [1]. However, Grant and Kalman [2] showed that the main mechanism controlling the grinding process in SAG mills is impact.

Comminution is the most energy-intensive part of the mineral recovery process [3]. However, it is believed that less than 5% of the energy supplied to comminution processes is used for particle breakage [4]; the rest is wasted through extra motions of the mill contents and collisions of the balls to the liners which not only waste energy but also damage the liners. SAG mills are lined with liners (lifters) to improve their efficiency. Lifters are an essential element in a grinding mill, whose purpose is to prevent slippage of the mill charge [5]. Without the lifters, the SAG mill cannot lift the charge high enough to promote impact breakage, and energy is consumed without substantial breakage. Liner life is the main factor influencing an industrial SAG mill's availability, which is an issue of significant economic importance, because one of the main reasons of mill down time is the time required to replace the worn or broken liners.

In this research, the equations governing the motion of grinding media inside a SAG mill and the impact force and stress due to the impact of steel balls onto the liners have been derived. The effect of the different design parameters have also been studied in order to optimize the performance of a SAG mill and avoid severe impacts causing liner breakage.

2- The Developed Model

In this study, a steel ball's motion, from the instant it rides the liner at one side of the mill until the time it falls down onto the liners on the other side, is divided into three phases. As shown in Figure 2, in the first phase, the steel ball is lifted by a lifter to a maximum height, described by the separation angle α , which the ball's position makes with the horizontal axis. At the instant that the ball is separated from the lifter, the normal force exerted by the wall N_w vanishes and the equations of the ball's motion in the x' and y' directions, shown in Figure 2, for a mill of radius R , angular velocity ω , and liner inclination angle φ , can be written as:

$$x' \text{- direction} \Rightarrow \mu.N - mg.\sin(\varphi + \alpha) = -mR\omega^2.\cos\varphi \quad (1-a)$$

$$y' \text{- direction} \Rightarrow N - mg.\cos(\varphi + \alpha) = mR\omega^2.\sin\varphi \quad (1-b)$$

The normal force N acting by the liner face on the ball is derived from the motion equation in the y' direction:

$$N = mg.\cos(\varphi + \alpha) + mR\omega^2.\sin\varphi \quad (2)$$

Substituting N into the motion equation in the x' direction and simplifying leads to:

$$\frac{R\omega^2}{g} [\mu.\sin\varphi + \cos\varphi] + [\mu.\cos(\varphi + \alpha) - \sin(\varphi + \alpha)] = 0 \quad (3-a)$$

Renaming the coefficient of friction as $\mu \equiv \tan\gamma = \frac{\sin\gamma}{\cos\gamma}$ and replacing into Eq. (3-a) gives:

$$\begin{aligned} & \frac{R\omega^2}{g} \left[\frac{\sin\varphi.\sin\gamma + \cos\varphi.\cos\gamma}{\cos\gamma} \right] + \left[\frac{\cos(\varphi + \alpha).\sin\gamma - \sin(\varphi + \alpha).\cos\gamma}{\cos\gamma} \right] = \\ & \frac{R\omega^2}{g} \left[\frac{-\sin\varphi.\sin(-\gamma) + \cos\varphi.\cos(-\gamma)}{\cos\gamma} \right] + \left[\frac{\cos(-(\varphi + \alpha)).\sin\gamma + \sin(-(\varphi + \alpha)).\cos\gamma}{\cos\gamma} \right] \quad (3-b) \\ & \Rightarrow \frac{R\omega^2}{g} \left[\frac{\cos(\varphi - \gamma)}{\cos\gamma} \right] + \left[\frac{\sin(\gamma - \varphi - \alpha)}{\cos\gamma} \right] = 0 \end{aligned}$$

which is further simplified to give the following equation for the separation angle:

$$\alpha = \arcsin \left[\frac{R \omega^2}{g} \cos(\varphi - \gamma) \right] + \gamma - \varphi. \quad (4)$$

During the second phase the ball moves down on the liner face, the details of which are illustrated in Figure 3 along with the forces acting on the steel ball. Depending on the design of the mill (R , ω , φ) and the coefficient of friction between the lifter face and the steel ball (μ), the motion of the steel ball down the liner face can either be a *pure rolling* motion or a *combined rolling and sliding* motion, each having their own set of governing equations. The following paragraph explains the procedure for establishing a criterion for the pure rolling motion.

In a pure rolling motion down the inclined face of the liner, the angular acceleration $\bar{\alpha}$ of the ball and the linear acceleration for pure rolling motion are as follows, respectively:

$$\sum M_c = I_c \bar{\alpha} \Rightarrow mg \sin(\alpha + \varphi) \times r_B = \left(\frac{5}{3} m r_B^2 \right) \bar{\alpha} \Rightarrow \bar{\alpha} = 0.6 \frac{g}{r_B} \sin(\alpha + \varphi) \quad (5-a)$$

$$\bar{a} = r_B \bar{\alpha} = 0.6 g \sin(\alpha + \varphi) \quad (5-b)$$

The force-acceleration equations for ball's motion in the t- and n- directions of Figure 3 in a pure rolling motion are:

$$t\text{-direction} \Rightarrow N \cdot \cos(\varphi) - mg \cos(\alpha) = 0 \Rightarrow N = mg \frac{\cos(\alpha)}{\cos(\varphi)} \quad (6-a)$$

$$n\text{-direction} \Rightarrow -F_\mu \cdot \cos(\varphi) + mg \sin(\alpha) = m \bar{a} = 0.6 mg \sin(\alpha + \varphi) \quad (6-b)$$

Simplifying Eq. (6-b) gives the friction force, which cannot exceed its maximum value:

$$F_\mu = mg \frac{\sin(\alpha) - 0.6 \sin(\alpha + \varphi)}{\cos(\varphi)} \leq \mu_s N \quad (7)$$

Finally, substituting the normal force N for pure rolling motion from Eq. (6-a) results in the criterion for pure rolling:

$$\mu_s \geq \frac{\sin(\alpha) - 0.6 \sin(\alpha + \varphi)}{\cos(\alpha)} \quad (8)$$

For a given mill with a fixed radius (R) and a constant angular velocity (ω), the pure rolling criterion can be restated by defining a function $F(\mu, \varphi)$ as follows:

$$F(\mu, \varphi) = \mu_s - \frac{\sin(\alpha) - 0.6 \sin(\alpha + \varphi)}{\cos(\alpha)} \quad (9)$$

Using the above notation, the pure rolling criterion can be stated as:

$$F(\mu, \varphi) \geq 0 \quad (10)$$

A program has been written in MATLAB to compute the pure rolling criterion for any specific SAG mill using Eq. (10).

In order to obtain realistic numerical solutions, the dimensions and operating conditions of a full-scale 32×16 ft (9.75×4.88 m) industrial SAG mill have been used. The above-mentioned SAG mill, the dimensions and operating conditions of which have been summarized in Table 1 [6], operates in Sarcheshmeh Copper Mining Complex in Rafsanjan, Iran.

Figure 4 shows the pure rolling criterion for different lifter angles (φ) and coefficients of friction (μ), for the Sarcheshmeh SAG mill, in which positive values in the curves resemble pure rolling motion down the liner face and negative values resemble combined rolling and sliding motion.

The results indicate that for the mill being studied, for liner angles greater than 15° the motion down the liner face is pure rolling, regardless of the coefficient of friction. For smaller liner angles, however, in the case of small coefficients of friction the motion of the steel balls is of combined rolling and sliding and for larger coefficients of friction it is pure rolling. For

rectangular liner profiles ($\varphi=0^\circ$), the motion is always combined rolling and sliding, regardless of the coefficient of friction.

If the ball moves down in a *pure rolling* motion, its angular velocity during the motion and linear velocity at the end of the path can be derived from the work and energy relation:

$$\frac{1}{2} \left(\frac{5}{3} m r_B^2 \right) \omega'^2 = mg \left(\frac{H}{\cos \varphi} - r_B \right) \sin(\alpha + \varphi) \Rightarrow \omega' = \sqrt{1.2 \frac{g}{r_B^2} \left[\frac{H}{\cos \varphi} - r_B \right] \sin(\alpha + \varphi)} \quad (10)$$

$$V = r_B \omega' = \sqrt{1.2 g \left(\frac{H}{\cos(\varphi)} - r_B \right) \sin(\alpha + \varphi)} \quad (12)$$

where H and r_B are the lifter height and the ball radius, respectively, and ω' is the angular velocity of the ball rolling down the liner face. Either the linear acceleration-displacement equation or the angular acceleration-velocity relation can be used to calculate the time it takes for the ball to roll down the liner face and reach the end of the path:

$$L = \frac{1}{2} \bar{a} t^2 \quad \text{or} \quad \omega' = \bar{\alpha} t \Rightarrow t = \sqrt{\frac{\left[\frac{H}{\cos(\varphi)} - r_B \right]}{0.3 g \sin(\alpha + \varphi)}} \quad (13)$$

In a *combined rolling and sliding* motion, the linear acceleration of the ball in the n -direction is obtained from the equation of motion in that direction:

$$\sum F_n = m \bar{a}_n \Rightarrow \bar{a} = g \sin(\alpha) - \mu_k g \frac{\cos(\alpha)}{\cos(\varphi)} \quad (14)$$

The angular acceleration of the ball $\bar{\alpha}$ can also be calculated as:

$$\sum M_G = I_G \bar{\alpha} \Rightarrow \mu_k m g \frac{\cos(\alpha)}{\cos(\varphi)} \times r_B = \frac{2}{3} m r_B^2 \bar{\alpha} \Rightarrow \bar{\alpha} = 1.5 \mu_k \frac{g}{r_B} \frac{\cos(\alpha)}{\cos(\varphi)} \quad (15)$$

The time it takes for the ball to reach the end of the path is:

$$L = \frac{H}{\cos(\varphi)} - r_B = \frac{1}{2} \bar{a} t^2 \Rightarrow t = \sqrt{\frac{2 \left[\frac{H}{\cos(\varphi)} - r_B \right]}{g \sin(\alpha) - \mu_k g \frac{\cos(\alpha)}{\cos(\varphi)}}} \quad (16)$$

Finally, the ball's velocity at the end of the path in a combined rolling and sliding motion is:

$$V = \bar{a} t \Rightarrow V = \sqrt{2g \left(\frac{H}{\cos(\varphi)} - r_B \right) \times \left(\sin(\alpha) - \mu_k \frac{\cos(\alpha)}{\cos(\varphi)} \right)} \quad (17)$$

The third phase of the motion, as shown in Figure 5, is a projectile motion in which the steel ball leaves the tip of the lifter with an initial velocity V_o and collides with a liner on the opposite side with velocity V_t . For simplicity, the calculations of the third phase of the motion are performed in Cartesian coordinates (x - and y - coordinates of Figure 3). From the geometry, the components of the initial position vector are:

$$x = R \cos(\alpha) - \frac{H}{\cos(\varphi)} \cos(\alpha + \varphi) \quad (18-a)$$

$$y = R \sin(\alpha) - \frac{H}{\cos(\varphi)} \sin(\alpha + \varphi) \quad (18-b)$$

Based on geometrical calculations, the x - and y - components of the initial velocity vector due to the mill's rotation (acting in the t -direction) are:

$$u_1 = -\omega \sqrt{R^2 + \frac{H^2}{\cos^2(\varphi)} - 2RH} \times \sin(\alpha - \psi) \quad (19-a)$$

$$v_1 = -\omega \sqrt{R^2 + \frac{H^2}{\cos^2(\phi)} - 2RH} \times \cos(\alpha - \psi) \quad (19-b)$$

in which the angle ψ , as shown in Figure 3, is determined as:

$$\psi = \text{Arc tan} \left[\frac{H \cdot \tan(\phi)}{R - H} \right] \quad (20)$$

There exists another velocity component for the ball, in the direction of the ball's motion down the face of the liner, denoted as the n' -direction. This velocity vector depends on the nature of the motion, whether it is pure rolling or it's rolling with slippage. The x- and y-components of this velocity vector, for pure rolling, are:

$$u_2 = -\sqrt{1.2g \left(\frac{H}{\cos(\phi)} - r_B \right)} \sin(\alpha + \phi) \times \cos(\alpha + \phi) \quad (21-a)$$

$$v_2 = -\sqrt{1.2g \left(\frac{H}{\cos(\phi)} - r_B \right)} \sin(\alpha + \phi) \times \sin(\alpha + \phi) \quad (21-b)$$

For rolling with sliding, they are:

$$u_2 = -\sqrt{2g \left(\frac{H}{\cos(\phi)} - r_B \right) \times \left(\sin(\alpha) - \mu_k \frac{\cos(\alpha)}{\cos(\phi)} \right)} \times \cos(\alpha + \phi) \quad (22-a)$$

$$v_2 = -\sqrt{2g \left(\frac{H}{\cos(\phi)} - r_B \right) \times \left(\sin(\alpha) - \mu_k \frac{\cos(\alpha)}{\cos(\phi)} \right)} \times \sin(\alpha + \phi) \quad (22-b)$$

From the moment the ball leaves the lifter and during its projectile motion, the only force acting on it is its weight. The initial velocity of the ball at the start of its projectile motion is:

$$U(0) = u_1 + u_2 \quad (23)$$

$$V(0) = v_1 + v_2 \quad (24)$$

Therefore, the components of the ball's velocity and the position of the ball at an instant t after the start of the motion can easily be calculated:

$$U(t) = u_1 + u_2 = \text{const.} \quad (25)$$

$$V(t) = -gt + (v_1 + v_2) \quad (26)$$

The position of the ball at the instant t is:

$$X(t) = (u_1 + u_2)t + x \quad (27)$$

$$Y(t) = -\frac{1}{2}gt^2 + (v_1 + v_2)t + y \quad (28)$$

For the ball to hit the lifter Eqs. (27) and (28) should be combined such that:

$$X^2 + Y^2 = R^2 \quad (29)$$

resulting in:

$$\left[(u_1 + u_2)t + x \right]^2 + \left[-\frac{1}{2}gt^2 + (v_1 + v_2)t + y \right]^2 = R^2 \quad (30)$$

Eq. (30) is then solved numerically for t , which in turn, can be used to give the angle β which dominates the position of the ball's impact, as shown in Figure 5:

$$\beta = \arctan \left(\frac{Y}{X} \right) \times \frac{180}{\pi} \quad (\text{deg}) \quad (31)$$

The impact velocity V_t and the impact angle θ can also be determined as follows:

$$|\vec{V}_{\text{impact}}| = \sqrt{U^2 + V^2} \quad (32)$$

$$\theta = \arctan\left(\frac{V(t)}{U(t)}\right) \times \frac{180}{\pi} \quad (\text{deg}) \quad (33)$$

The effective impact angle, defined as the acute angle between the liner face surface and the velocity vector, can easily be figured out from the geometry of Figure 3 as:

$$\theta' = \beta - \theta \quad (34)$$

The normal or tangent components of the impact velocity vector can be easily calculated by multiplying $|\vec{V}_{\text{impact}}|$ in $\text{Sin } \theta'$ or $\text{Cos } \theta'$, respectively.

The deformation and stress distribution in the contact region, due to the elastic impact of spheres can be satisfactorily accounted for using the stress analysis for curved, contacting solid bodies presented by Hertz [7]. We assume that motion occurs only along the line joining the mass centers, the deformations due to impact are local, and deformations near the mass centers are negligible. The Hertz's Law of Contact, in the contact of two spheres, combined with Newton's second law form [8]:

$$F = K_2 \delta^{\frac{3}{2}} = -\frac{\ddot{\delta}}{\left(\frac{1}{m_1} + \frac{1}{m_2}\right)} = -\frac{\ddot{\delta}}{K_1} \quad (35)$$

in which δ (approach) is the distance through which the centers of the two spheres approach due to the local compression and:

$$K_1 = \frac{m_1 + m_2}{m_1 m_2} = \frac{1}{m_1} + \frac{1}{m_2} \quad (36-a)$$

$$K_2 = \frac{4}{3(k_1 + k_2)} \sqrt{\frac{R_1 R_2}{R_1 + R_2}} = \frac{4}{3(k_1 + k_2)} \sqrt{\frac{1}{R_1} + \frac{1}{R_2}} \quad (36-b)$$

where

$$k_1 = \frac{1 - \nu_1^2}{E_1} \quad (37-a)$$

$$k_2 = \frac{1 - \nu_2^2}{E_2} \quad (37-b)$$

By integrating Eq. (35) and applying the following initial condition:

$$\dot{\delta} = v_0 \quad \text{at } \delta = 0 \quad (38)$$

we have:

$$\frac{1}{2}(\dot{\delta}^2 - v_0^2) = -\frac{2}{5} K_1 K_2 \delta^{\frac{5}{2}} \Rightarrow \dot{\delta}^2 = v_0^2 - \frac{4}{5} K_1 K_2 \delta^{\frac{5}{2}} \quad (39)$$

Maximum compression δ_m , which occurs at the instant of zero relative velocity ($\dot{\delta} = 0$), is easily derived from Eq. (39):

$$\delta_m = \left[\frac{5v_0^2}{4K_1 K_2} \right]^{\frac{2}{5}} \quad (40)$$

Assuming that the time history of deformation is symmetric about the moment at which the maximum deformation occurs and since:

$$dt = \frac{d\delta}{\dot{\delta}} \quad (41)$$

the impact duration can be expressed as:

$$T_H = \int_0^{T_H} dt = 2 \int_0^{\frac{T_H}{2}} dt = 2 \int_0^{\delta_m} \frac{d\delta}{\delta} \quad (42)$$

Replacing from Eq. (39) gives:

$$T_H = 2 \int_0^{\delta_m} \frac{d\delta}{\sqrt{v_0^2 - \frac{4}{5} K_1 K_2 \delta^2}} \quad (43)$$

Eq. (43) can, in turn, be simplified to the form of:

$$T_H = \frac{2}{v_0} \int_0^{\delta_m} \frac{d\delta}{\sqrt{1 - (\delta/\delta_m)^2}} \quad (44)$$

After the substitution $Z = \frac{\delta}{\delta_m}$ and $d\delta = \delta_m dZ$, we have:

$$T_H = \frac{2\delta_m}{v_0} \int_0^1 \frac{dZ}{\sqrt{1 - Z^2}} \quad (45)$$

Finally, by using the transformations $u = 1 - Z^2$, we arrive at:

$$T_H = \frac{4\delta_m}{5v_0} \int_0^1 u^{-\frac{1}{2}} (1-u)^{-\frac{3}{5}} du \quad (46)$$

which is an Eulerian Integral of the first kind [9], the solution of which is a Beta Function [10]:

$$T_H = \frac{4\delta_m}{5v_0} \beta\left(\frac{1}{2}, \frac{2}{5}\right) \Rightarrow T_H \approx \frac{2.943275}{(v_0)^{\frac{1}{5}}} \left[\frac{5}{4K_1 K_2} \right]^{\frac{2}{5}} \quad (47)$$

For the impact of a sphere of mass m and radius R_1 with a heavy plate ($R_2 \rightarrow \infty$, $m_2 \rightarrow \infty$):

$$K_1 \rightarrow \frac{1}{m}, \quad K_2 = \frac{4\sqrt{R_1}}{3(k_1 + k_2)} \quad (48)$$

and, therefore, Eq. (47) simplifies to the form:

$$T_H = \frac{2.943275}{v_0^{\frac{1}{5}}} \left[\frac{5m}{4K_2} \right]^{\frac{2}{5}} = \frac{2.868266}{v_0^{\frac{1}{5}}} \left[\frac{m(k_1 + k_2)}{\sqrt{R_1}} \right]^{\frac{2}{5}} \quad (49)$$

If the plate and the sphere be made of the same material, Eq. (49) is simplified in terms of material properties to give:

$$T_H = 3.356115 \left[\rho \frac{1 - \nu^2}{E} \right]^{\frac{2}{5}} \frac{d}{(v_0)^{\frac{1}{5}}} \quad (50)$$

The elastic and inertial properties of typical steel balls and SAG mills liners are:

$$\begin{aligned} E &= 200 \text{ GPa} \\ \nu &= 0.29 \\ e &= 0.49 \\ \rho &= 7800 \frac{\text{Kg}}{\text{m}^3} \end{aligned} \quad (51)$$

Replacing the values for the elastic and inertial properties results in:

$$T_H = 3.52372 \times 10^{-3} d (v_0)^{-\frac{1}{5}} \quad (52)$$

The impacts studied here are impacts in which permanent plastic deformations do not occur and the deformations during the impact period are elastic. In this kind of impact, the time history of the magnitude of impact forces can be estimated as a triangle [11], as in Figure 6. The impact period consists of a deformation period (from 0 to t_0) and a restoration period (from t_0 to t). The short period of deformation takes place right after the initial contact of the two bodies and continues until the contact area stops increasing. Then, the period of restoration starts in which the contact area decreases until it ultimately returns back to zero. The coefficient of restitution is defined as:

$$e = \frac{\int_{t_0}^t F_r dt}{\int_0^{t_0} F_d dt} = \frac{(v_2')_n - (v_1')_n}{(v_1)_n - (v_2)_n} \quad (53)$$

in which F_d and F_r are the normal forces during deformation and restoration, respectively and $(v_i)_n$ and $(v_i')_n$ are the n-components of the velocity of body i ($i=1,2$) before and after the impact, respectively.

If one of the impacting bodies (e.g. body 2) is stationary before impact and remains stationary after the impact ($v_2' = v_2 = 0$), as in the case of the wall in the collision of a ball against a rigid stationary wall, the coefficient of restitution reduces to:

$$e = -\frac{(v_1')_n}{(v_1)_n} = -\frac{(v_1')_n}{v_0} \quad (54)$$

The Linear Momentum relation in the direction of impact for one of the bodies is used to calculate the maximum impact force P_{max} :

$$\int_{t_1}^{t_2} (\sum F) dt = \frac{1}{2} P_{max} T_H = \Delta G = m [(v_1')_n - (v_0)] = m v_0 (1+e) \quad (55)$$

Replacing the coefficient of restitution and simplifying results in the maximum impact force:

$$P_{max} = \frac{-2m_1 (v_1)_n (1+e)}{T_H} = \frac{-2 \left(\rho \pi \frac{d^3}{6} \right) (v_1)_n (1+e)}{T_H} \quad (56)$$

Substituting the impact time from Eq. (52) and simplifying for the impact of a steel ball of diameter d gives:

$$P_{max} = -3.45400 \times 10^6 \times [(v_1)_n]^{1.2} d^2 \quad (57)$$

If the two contacting bodies are made of the same material, following the procedure used in [8] and fully explained in [12], the radius b of the contact circle at the point of impact can be satisfactorily approximated by:

$$b = 0.90 \sqrt[3]{P_{max} \Delta} \quad (58)$$

in which $\Delta = 2R \left(\frac{1-\nu^2}{E} \right)$. The maximum principal stress σ_{max} is:

$$\sigma_{max} = -0.65 \left(\frac{b}{\Delta} \right) \quad (59)$$

The maximum shear stress τ_{max} is:

$$\tau_{max} = 0.21 \left(\frac{b}{\Delta} \right) \quad (60)$$

The maximum octahedral shear stress $\tau_{oct(max)}$ is:

$$\tau_{oct(max)} = 0.196 \left(\frac{b}{\Delta} \right) \quad (61)$$

The maximum shear stresses τ_{max} and $\tau_{oct(max)}$ occur at the depth z_s from the free surface:

$$z_s = 0.47b \quad (62)$$

The maximum approach δ_m is:

$$\delta_m = 2.40 \left(\frac{P_{max} \Delta}{\pi b R} \right) \quad (63)$$

3- Results and Discussion

The most important step in optimization processes is to identify the main design variable. One of the main variables characterizing the motion of steel balls in a SAG mill is the **separation angle α** , shown schematically in Figure 2. As denoted by Eq. (4), the most important design variables and operating conditions affecting the separation angle in an industrial SAG mill are:

- Mill radius R ;
- Mill rotational velocity ω ;
- Liner inclination angle φ ;
- Liner surface coefficient of static friction μ_s .

The main characteristics dominating the impact of steel balls onto SAG mill shell liners are the **impact position (X and Y or the angle β)**, **impact velocity (V_{impact})**, **impact angle (θ')**, and **various impact stresses**. As demonstrated by Eqs. (18)-(19), (21)-(22) and (25)-(28) the design variables and operating conditions of an industrial SAG mill affecting the impact position, the impact velocity, the impact angle and various impact stress values are:

- Mill radius R ;
- Mill rotational velocity ω ;
- Liner inclination angle φ ;
- Liner surface coefficient of static friction μ_s ;
- Lifter height H ;
- Steel ball size (radius) r_B ;
- Liner surface coefficient of kinetic friction μ_k (in the case of *combined rolling and sliding* motion down the liner face).

One approach to optimizing the motion of steel balls in a SAG mill is to separately measure the effect of each single one of the afore-mentioned design variables on the characteristics of the motion of steel balls and their impact onto SAG mill shell liners, while keeping all the other variables constant, in order to determine the net effect of each variable on the operating condition.

A second approach for determining the overall effect of these design variables on the operating condition is to measure the effect of a number of the design variables on the characteristics of steel ball motion and its impact onto SAG mill shell liners simultaneously, while keeping the rest of the design variables constant.

As a demonstrative case study, the second approach has been followed in this paper and the effect of two variables, namely the lifter height (H) and the coefficient of friction (μ) on some of the main impact characteristics have been simultaneously investigated for the Sarcheshmeh SAG mill.

In order to determine the effect of the lifter height and the coefficient of friction on the impact position (β), the impact positions corresponding to various lifter heights were plotted

against the coefficient of friction (μ) for different combinations of lifter face angle (ϕ) and steel ball diameter (d_B). Figure 7 illustrates the impact position for a lifter face inclination angle of 15° and a 2-inch diameter steel ball. Results indicate that the impact position tends to move upward (the angle β decreases) as the lifter height or the coefficient of friction increases. This is the direct result of the upward move of the separation point due to the increase in the time it takes for the ball to roll or slide down the lifter face in higher lifters or for lifters which have a higher coefficient of friction.

To investigate the effect of the coefficient of friction and the lifter height on the maximum impact force (P_{max}), the maximum impact forces corresponding to various lifter heights were plotted against the coefficient of friction (μ) for various steel ball diameters (d_B) and lifter face angles (ϕ). Figure 8 illustrates the results for a lifter face inclination angle of 15° and a steel ball diameter of 2 inches. As shown in Figure 8, the maximum impact force decreases as the lifter height or the coefficient of friction increase.

In order to demonstrate the effect of the coefficient of friction and the lifter height on the maximum principal stress (σ_{max}), the curves corresponding to the maximum principal stresses due to different lifter heights were plotted against the coefficient of friction (μ) for various combinations of steel ball size (d_B) and lifter face inclination angle (ϕ). Figure 9 shows the maximum principal stress for a lifter face angle of 15° and a 2-inch diameter steel ball. Results indicate that the absolute value of the maximum principal stress decreases as the lifter height or the coefficient of friction increase.

4- Conclusion

In this research, the equations governing the motion of steel balls in an industrial SAG mill and their impact onto shell liners were derived in full details and were used in order to determine the effective parameters for optimizing the working conditions of the mill and to avoid severe impacts which lead to the breakage of SAG mill shell liners. Based on the aforementioned equations, the most important design variables governing the intensity of the impacts are lifter height H , the working coefficient of friction μ , lifter face inclination angle ϕ , steel ball size r_B , mill rotational velocity ω , and mill size (radius) R . In order to optimize the operating conditions of a SAG mill and avoid severe impacts to its shell liners, which lead to the breakage of liners, the effect of these parameters need to be studied.

As a case study, the effect of the lifter height (H) and the coefficient of friction (μ) on some of the main impact characteristics were simultaneously investigated for an industrial SAG mill installed at Sarcheshmeh Copper Complex. It was shown that the lifter height (H) and the coefficient of friction (μ) have a significant effect on impact parameters and therefore on liner damage. Based on the results of the case study, as the lifter height or the coefficient of friction increases the impact position tends to move upward (the angle β decreases).

Moreover, considering the maximum impact force it was shown that the maximum impact force decreases as the lifter height or the coefficient of friction increase reducing the impact severity. Results indicated that the absolute value of the maximum principal stress decreases as the lifter height or the coefficient of friction increase.

Therefore, this study and its results emphasize the importance of determining the most important design variables and working conditions which influence the motion of steel balls in a SAG mill. The effect of these design variables can be determined in an optimization process and the results can be used in order to change the design of SAG mills or their working conditions to avoid damaging impacts onto the liners and hence liner breakage due to severe impacts.

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Nomenclature

- \bar{a} : Linear acceleration of steel ball
 d_B : Steel ball diameter
 e : Coefficient of restitution
 F_μ : Friction force
 H : Lifter height
 m : Steel Ball mass
 N : Normal force acting on the ball from the lifter
 N_w : Normal force acting on the ball from the wall
 P_{\max} : Maximum impact force
 R : Mill radius
 r_B : Steel ball radius
 T_H : Impact duration
 Z_s : Depth at which maximum shear stresses occur

Greek symbols

- α : Separation angle
 $\bar{\alpha}$: Angular acceleration of steel ball
 β : Impact position with reference to the horizontal
 δ : Approach
 δ_m : Maximum approach
 φ : Lifter face angle
 μ_s : Static coefficient of friction
 μ_k : Kinetic coefficient of friction
 θ : Impact angle
 θ' : Effective impact angle
 σ : Principal stress
 τ : Shear stress
 τ_{oct} : Octahedral shear stress
 ω : Angular velocity of the mill
 ω' : Angular velocity of steel ball

Tables

Table 1 Sarcheshmeh SAG mill Specifications

Speed, rpm	10.52 rpm
Critical Speed, rpm	13.66 rpm
Speed, % of critical speed	77% V_c
Max Charge Filling, % of total mill volume	35%
Ball Filling, % of total mill volume	15%
Maximum Steel Ball Diameter	127 mm
Nominal Mill Length	4.878 m
Effective Grinding Length	4.420 m
Inside Shell Diameter	9.754 m
Aspect Ratio, diameter/length	2.2
Power Consumption	8.2 MW
Number of Liners	2 × 60
Liner Thickness	78 mm
Lifter Height (integral on the liner)	150 mm
Lifter Inclination Angle	15°

Figures

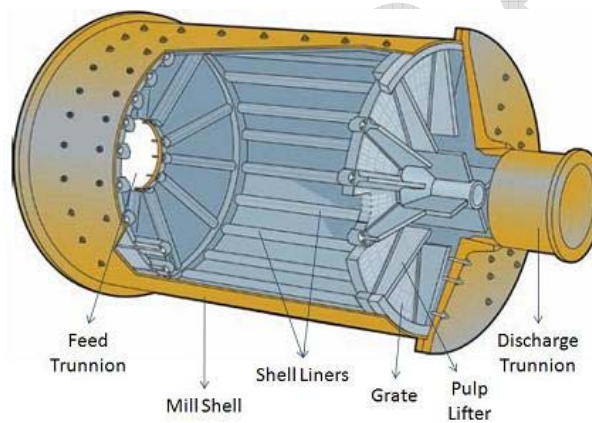


Figure 1 Schematic Design of a typical SAG mill

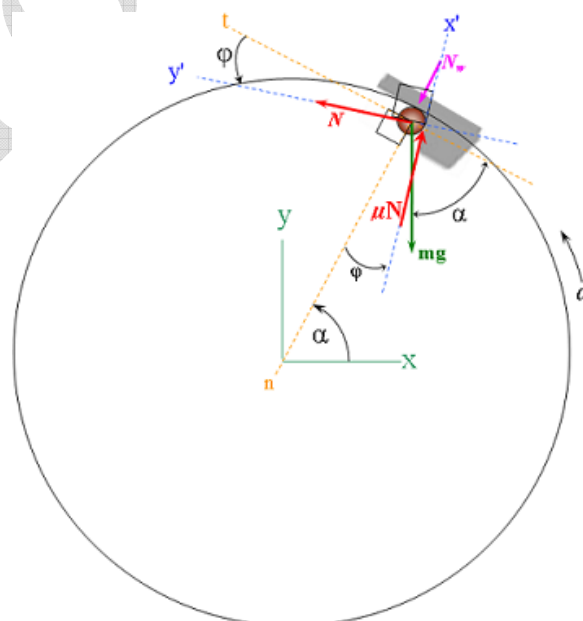


Figure 2 The first phase of the motion of a steel ball

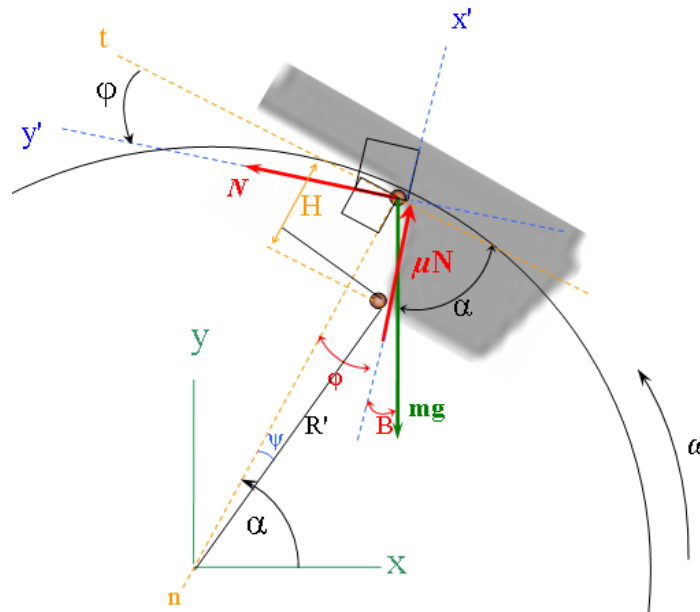


Figure 3 The second phase of the motion of a steel ball

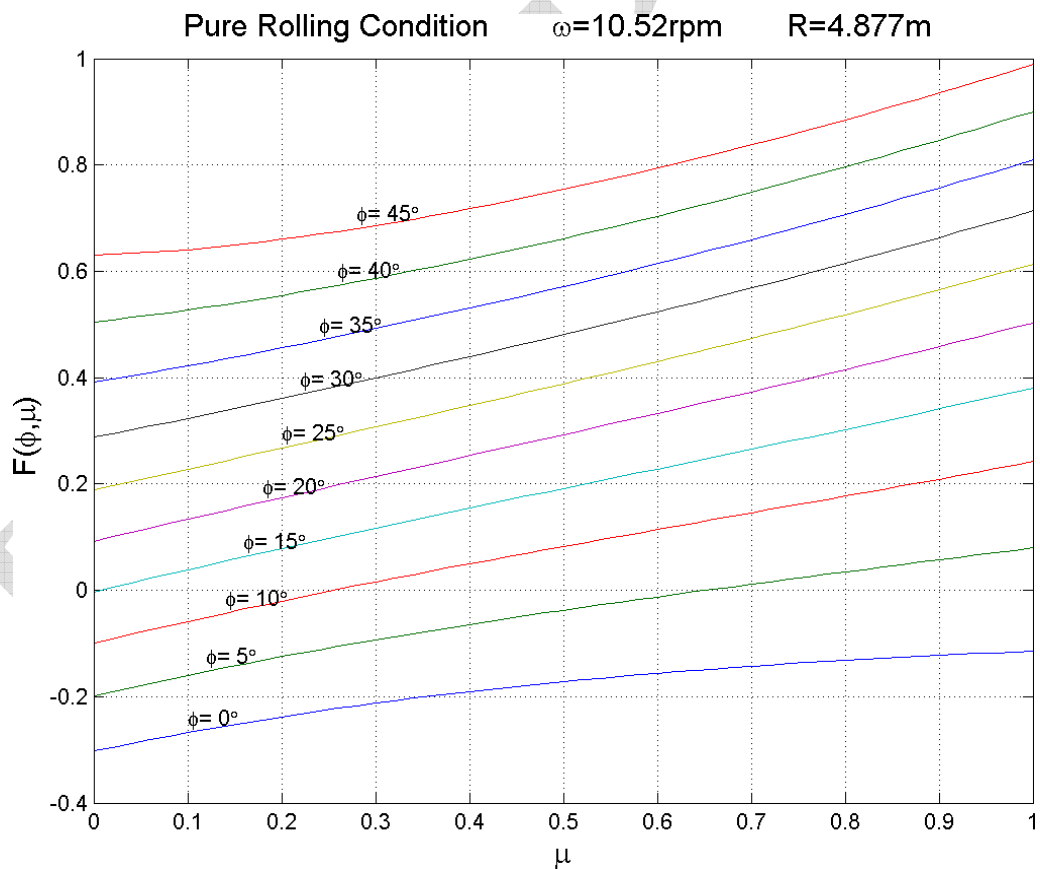


Figure 4 The pure rolling criterion

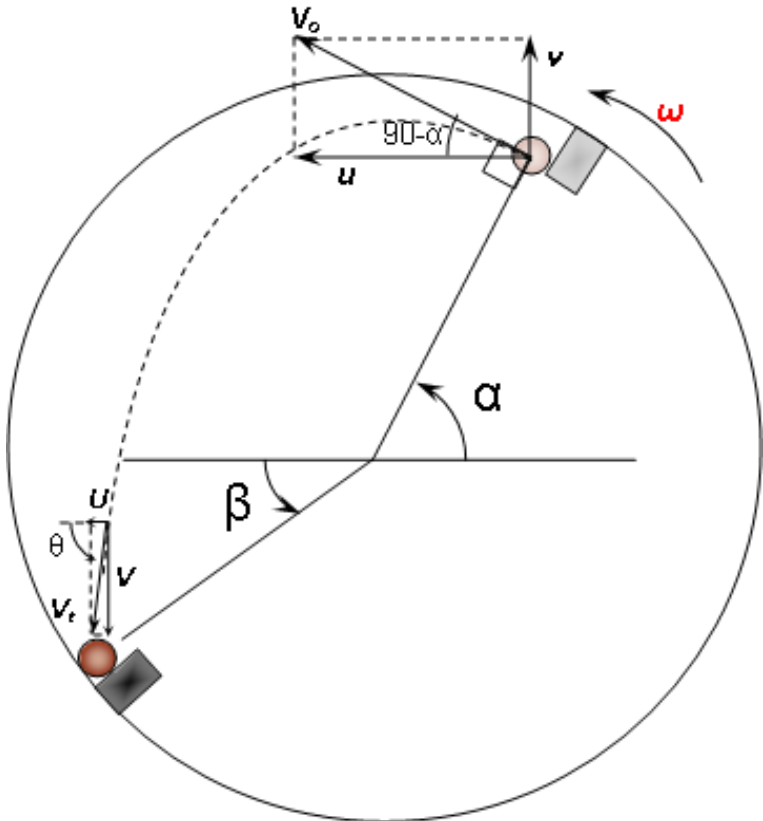


Figure 5 The third phase of a steel ball's motion

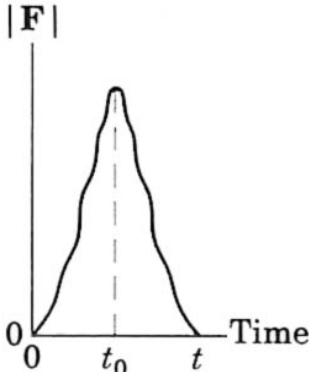


Figure 6 Time history of the impact force magnitude

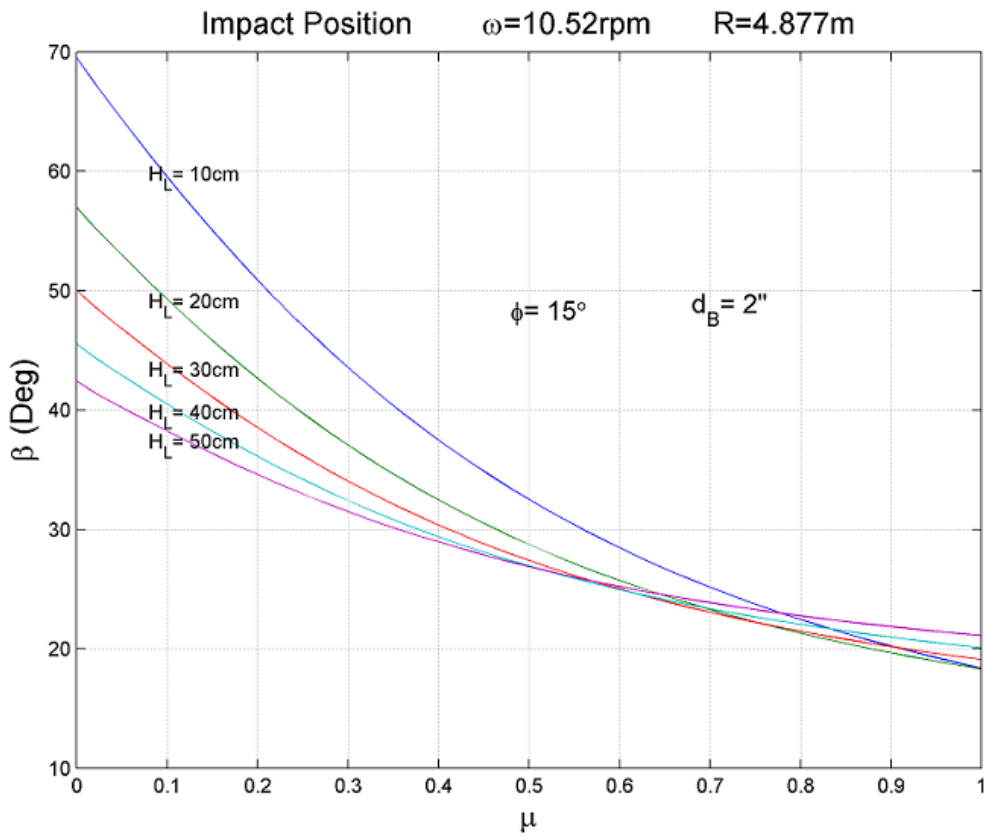


Figure 7 The effect of lifter height and the coefficient of friction on the impact position (Lifter face angle $\phi=15^\circ$, Ball diameter $d_B=2$ inches)

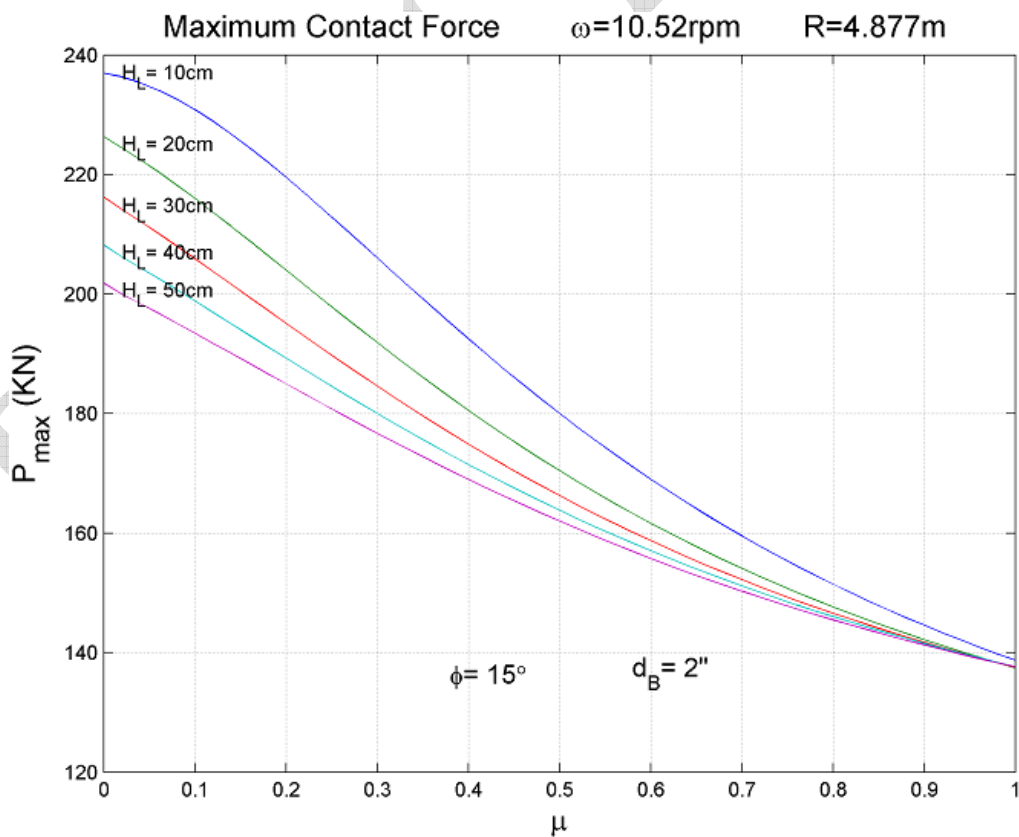


Figure 8 The effect of the coefficient of friction and lifter height on the maximum impact force (Lifter face angle $\phi=15^\circ$, Ball diameter $d_B=2$ inches)

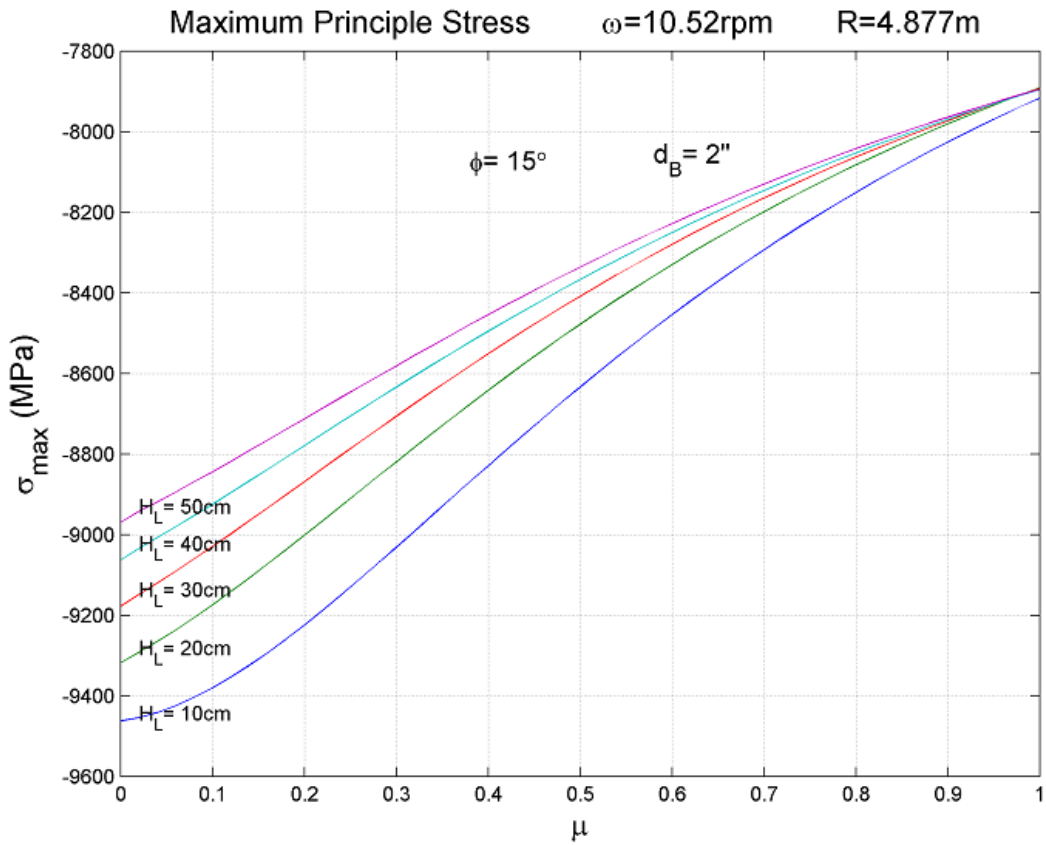


Figure 9 The effect of lifter height and the coefficient of friction on the maximum principal stress (Lifter face angle $\phi=15^\circ$, Ball diameter $d_B=2$ inches)

PROOF

(Shell Liners)

(SAG Mill)

μ H (Lifter) :

 .R ω r_B ϕ

 (μ) (H)

است.

Proof Read